

NEWTON SOLVER STABILIZATION FOR STOKES SOLVERS IN GEODYNAMIC PROBLEMS

Fraters M. R. T. (1), Bangerth W. (2), Thieulot C. (1), Glerum A. (1,3), Spakman W. (1,4) (1) DEPARTMENT OF EARTH SCIENCES, UTRECHT UNIVERSITY, NETHERLANDS, (2) DEPARTMENT OF MATHEMATICS, COLORADO STATE UNIVERSITY, U.S.A, (3) GFZ GERMAN RESEARCH CENTER FOR GEOSCIENCES, POTSDAM, GERMANY, (4) CENTRE OF EARTH EVOLUTION AND DYNAMICS (CEED), UNIVERSITY OF OSLO, NORWAY

ABSTRACT

- 1. We have implemented a Newton solver for ASPECT (Kronbichler et al., 2012)
- 2. We only use solver libraries to solve linear systems (no Trilinos NOX or PETSC SNES)
- 3. We have implemented line search and oversolving-prevention ourselves
- 4. The Jacobian is not always Symmetric Positive Definite (SPD)
- 5. We force the Jacobian to be SPD in a cheap and optimal way
- 6. This allows for more complex rheologies to be used with a Newton solver

PROBLEM STATEMENT

We are interested solving the Stokes equations:

$$\begin{split} -\nabla\cdot\left[2\eta\left(\dot{\boldsymbol{\epsilon}}(\boldsymbol{u})-\frac{1}{3}(\nabla\cdot\boldsymbol{u})\boldsymbol{1}\right)\right]+\nabla p&=\rho\boldsymbol{g}\qquad\qquad\text{in }\Omega,\\ \nabla\cdot\left(\rho\boldsymbol{u}\right)&=0\qquad\qquad\qquad\text{in }\Omega, \end{split}$$

The weak form of the Newton linearisation:

$$\begin{pmatrix} J_k^{uu} & J_k^{up} \\ J_k^{pu} & 0 \end{pmatrix} \begin{pmatrix} \delta U_k \\ \delta \bar{P}_k \end{pmatrix} = \begin{pmatrix} F_k^u \\ F_k^p \end{pmatrix}$$

where for the incompressible case the Jacobian elements are (ignoring the pressure scaling):

$$\begin{split} \left(J_k^{uu}\right)_{ij} &= \left(A_k\right)_{ij} + \left(\varepsilon(\varphi_i^u), 2\left(\frac{\partial\eta(\varepsilon(\boldsymbol{u}_k), p_k)}{\partial\varepsilon} : \varepsilon(\varphi_j^u)\right)\varepsilon(\boldsymbol{u}_k)\right), \\ \left(J_k^{up}\right)_{ij} &= B_{ij}^T + \left(\varepsilon(\varphi_i^u), 2\left(\frac{\partial\eta(\varepsilon(\boldsymbol{u}_k), p_k)}{\partial p}\varphi_j^p\right)\varepsilon(\boldsymbol{u}_k)\right), \\ \left(J_k^{pu}\right)_{ij} &= B_{ij}. \end{split}$$

In general J^{uu} neither symmetric, nor positive definite, which can be very bad for solvers.

PRELIMINARY CONCLUSIONS

The Newton solver without stabilisation is:

- fast;
- prone to numerical breakdowns;
- very sensitive to the precise tweaking of parameters such as minimum linear tolerance, line search iterations, etc.

Restoring Symmetry

Approx. J^{uu} with the eq. below where $E^{sym} = \frac{1}{2} \left(E_{mnpq} + E_{pqmn} \right)$ and $E(\varepsilon(u_k))_{mnpq} = \left[2\varepsilon(u)_{mn} \frac{\partial \eta(\varepsilon(u), p)}{\partial \varepsilon_{pq}} \right]$:

$$(J_k^{uu})_{ij} \approx (A_k)_{ij} + \left(\varepsilon(\varphi_i^u), \left(\frac{\partial\eta(\varepsilon(\boldsymbol{u}_k), p_k)}{\partial\varepsilon} : \varepsilon(\varphi_j^u)\right)\varepsilon(\boldsymbol{u}_k)\right)$$

= $(A_k)_{ij} + \left(\varepsilon(\varphi_i^u), E^{\text{sym}}(\varepsilon(\boldsymbol{u}_k))\varepsilon(\varphi_j^u)\right)$

FORCING SYMMETRY AND POSITIVE DEFINITENESS

The positive definiteness of J^{uu} is determined by tensor H:

$$J^{uu} = \left(\varepsilon(\varphi_i^u), 2\eta(\varepsilon(u))\varepsilon(\varphi_j^u)\right) + \left(\varepsilon(\varphi_i^u), E^{\text{sym}}(\varepsilon(u))\varepsilon(\varphi_j^u)\right)$$
$$= \left(\varepsilon(\varphi_i^u), \underbrace{\left[2\eta(\varepsilon(u))I \otimes I + E^{\text{sym}}(\varepsilon(u))\right]}_{=:H}\varepsilon(\varphi_j^u)\right)$$

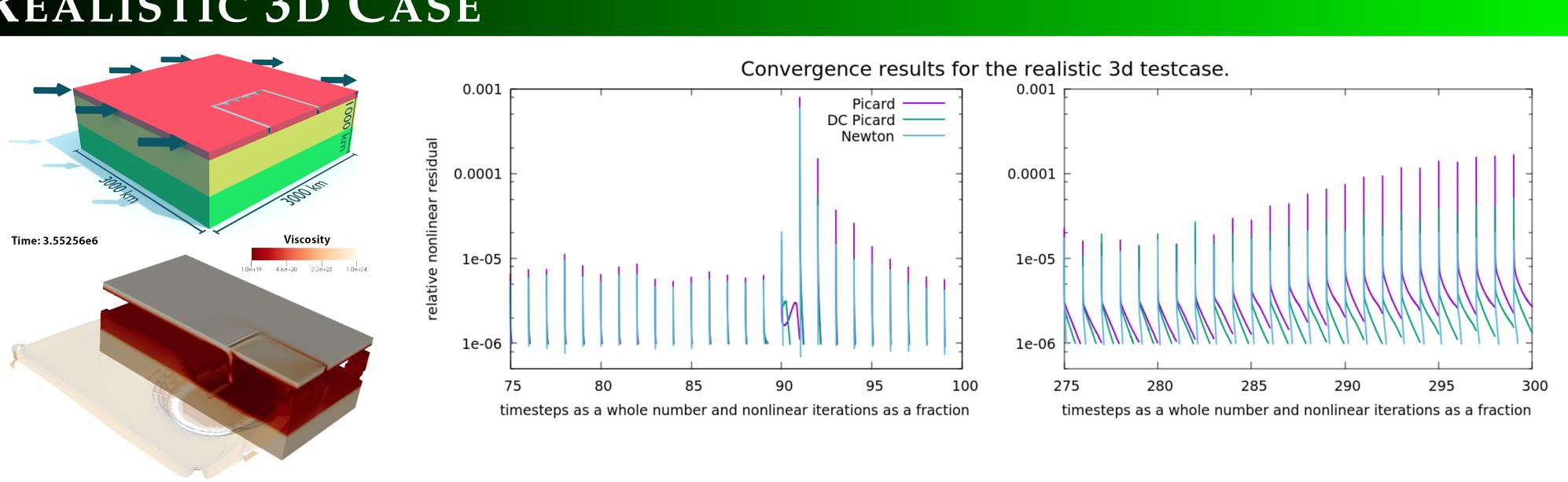
We can force tensor *H* to be SPD by scaling E^{sym} with a factor α :

$$H^{\mathrm{spd}} = 2\eta(\varepsilon(u))I \otimes I + \alpha E$$

with $0 < \alpha \leq 1$. If $\alpha = 0 J^{uu}$ will always be SPD, but to have optimal convergence we want α to be as large as possible (max one). It can be proven (Fraters et al., in prep) that the optimal value for α is (where $a = \varepsilon(u)$ and $b = \frac{\partial \eta(\varepsilon(u), p)}{\partial \varepsilon}$):

$$\alpha = \begin{cases} 1 & \text{if } \left[1 - \frac{1}{\|a\|}\right] \\ \frac{2\eta(\varepsilon(\boldsymbol{u}))}{\left[1 - \frac{b:a}{\|a\|}\right]^2 \|a\| \|b\|} & \text{otherwise.} \end{cases}$$

REALISTIC 3D CASE



The Newton solver with stabilisation is:

- faster than Picard, slower than without stabilisation;
- almost immune to numerical breakdowns;
- very insensitive to tweaking of these parameters.

 $+\left(\varepsilon(\varphi_{j}^{u}),\left(rac{\partial\eta(\varepsilon(\boldsymbol{u}_{k}),p_{k})}{\partial\varepsilon}:\varepsilon(\varphi_{i}^{u})
ight)\varepsilon(\boldsymbol{u}_{k})
ight)$

 ${\mathbb E}^{\mathrm{sym}}(arepsilon(u))$

 $\frac{b:a}{\|a\| \|b}$

$$\mathbf{x}_{\mathrm{T}} \Big]^2 \|a\| \|b\| < 2\eta(\varepsilon(\boldsymbol{u}))$$

BENCHMARKING

We show here results of a setup based on Spiegelman et al. (2016), on a mesh that has been refined either 4 (64x16 cells) or 8 times (1024x256 cells) uniformly.

